# Homework 4: Solutions to exercises not appearing in Pressley, also 2.3.2, 2.3.4, and 2.3.5 

Math 120A

- (2.3.2) We already know that the helix $\gamma(t)=(a \cos \theta, a \sin \theta, b \theta)$ with $a>0$ has curvature $\kappa=\frac{a}{a^{2}+b^{2}}$ and $\tau=\frac{b}{a^{2}+b^{2}}$. Since any two curves in $\mathbb{R}^{3}$ with the same nonzero curvature function and torsion functions are related by direct isometry, it suffices to show that for any pair of real numbers $\kappa>0$ and $\tau$, there is a pair $a>0$ and $b$ such that $\kappa=\frac{a}{a^{2}+b^{2}}$ and $\tau=\frac{b}{a^{2}+b^{2}}$. We see that $\frac{\tau}{\kappa}=\frac{b}{a}$, so we must have $b=\tau c$ and $a=\kappa c$ for some constant $c$. But then $\kappa=\frac{\kappa c}{\kappa^{2} c^{2}+\tau^{2} c^{2}}$, so $c=\frac{1}{\kappa^{2}+\tau^{2}}$. Ergo if $a=\frac{\kappa}{\kappa^{2}+\tau^{2}}$ and $b=\frac{\tau}{\tau^{2}+\kappa^{2}}$, the helix $\gamma(t)=(a \cos \theta, a \sin \theta, b \theta)$ has constant curvature and torsion $\kappa$ and $\tau$. We conclude that all curves with constant curvature and torsion are the images of circular helices under direct isometry.
- (2.3.4) We know $\gamma(t)$ is unit-speed and spherical. Without loss of generality, the center of the sphere is $\mathbf{0}$ (because if it isn't we can start by doing a translation). Therefore $\gamma(t) \cdot \gamma(t)=r^{2}$ where $r$ is the radius of the sphere. Differentiating gives $2 \dot{\gamma} \cdot \gamma=0$. SInce $\gamma$ is unit-speed, $\dot{\gamma}=\mathbf{t}$, so this says $\mathbf{t} \cdot \gamma=0$. Differentiating this second relationship again gives

$$
\begin{aligned}
\mathbf{t} \cdot \mathbf{t}+\kappa \mathbf{n} \cdot \gamma & =0 \\
1+\kappa \mathbf{n} \cdot \gamma & =0 \\
\mathbf{n} \cdot \gamma & =-\frac{1}{\kappa}
\end{aligned}
$$

We differentiate the last equality again:

$$
\begin{aligned}
(-\kappa \mathbf{t}+\tau \mathbf{b}) \cdot \gamma+\mathbf{n} \cdot \mathbf{t} & =\frac{\dot{\kappa}}{\kappa^{2}} \\
\tau \mathbf{b}) \cdot \gamma & =\frac{\dot{\kappa}}{\kappa^{2}} \\
\mathbf{b} \cdot \gamma & =\frac{\dot{\kappa}}{\kappa^{2} \tau}
\end{aligned}
$$

Notice this implies that $\gamma(t)=-\frac{1}{\kappa}+\frac{\dot{\kappa}}{\kappa^{2} \tau} \mathbf{n}$ for all $t$. We differentiate a final time:

$$
\begin{aligned}
-\tau \mathbf{n} \cdot \gamma+\mathbf{b} \cdot \mathbf{t} & =\frac{d}{d t}\left(\frac{\dot{\kappa}}{\kappa^{2} \tau}\right) \\
-\tau\left(-\frac{1}{\kappa}\right) & =\frac{d}{d t}\left(\frac{\dot{\kappa}}{\kappa^{2} \tau}\right) \\
\frac{\tau}{\kappa} & =\frac{d}{d t}\left(\frac{\dot{\kappa}}{\kappa^{2} \tau}\right)
\end{aligned}
$$

This completes the forward direction. Conversely, if this equation holds and $\rho=\frac{1}{\kappa}$, $\sigma=\frac{1}{\tau}$, then

$$
\begin{aligned}
\frac{d}{d t}\left(\rho^{2}+(\dot{\rho} \sigma)^{2}\right) & =2 \rho \dot{\rho}+2 \dot{\rho} \sigma \frac{d}{d t}(\dot{\rho} \sigma) \\
& =2 \dot{\rho}\left(\rho+\sigma \frac{d}{d t}(\dot{\rho} \sigma)\right)
\end{aligned}
$$

But $\rho+\sigma \frac{d}{d t}(\dot{\rho} \sigma)=\frac{1}{\kappa}+\frac{1}{\tau} \frac{d}{d t}\left(\frac{-\dot{k}}{\kappa^{2} \tau}\right)=0$ by equation (2.22). So we conclude that $\rho^{2}+(\dot{\rho} \sigma)^{2}$ is constant, say $r^{2}$. Now we know $\rho^{2}+(\dot{\rho} \sigma)^{2}=r^{2}$. So the curve $-\rho \mathbf{n}+(\dot{\rho} \sigma) \mathbf{b}$ lies on the sphere of radius $r$. Differentiating shows it has the same tangent vector as $\gamma$, so up to translation (which just changes the center of the sphere) it is the same curve. Ergo $\gamma$ is spherical. A computation shows the relationship holds for Viviani's Curve.

- (2.3.5) Let $\gamma(t)$ be our unit-speed curve and $P \mathbf{x}+\mathbf{a}$ be our direct isometry. Then if $\Gamma(t)=P \gamma(t)+\mathbf{a}$, we see that $\Gamma^{\prime}(t)=P \gamma^{\prime}(t)$, and since $P$ is orthogonal (and in particular, length-preserving), $\left\|P \gamma^{\prime}(t)\right\|=1$, so $\Gamma(t)$ is unit speed. Since both $\gamma$ and $\Gamma$ are unit speed, $\Gamma^{\prime}(t)=P \gamma^{\prime}(t)$ is exactly the statement that $\mathbf{T}=P \mathbf{t}$; similarly, $\Gamma^{\prime \prime}(t)=P \gamma^{\prime \prime}(t)$, so since $\left\|\gamma^{\prime \prime}(t)\right\|=\left\|P \gamma^{\prime \prime}(t)\right\|$, we see that $N=\frac{\Gamma^{\prime \prime}(t)}{\left\|\Gamma^{\prime \prime}(t)\right\|}=\frac{P \gamma^{\prime \prime}(t)}{\| \gamma^{\prime \prime}(t)} \|=P \mathbf{n}$. Then $P \mathbf{t} \times P \mathbf{n}=\operatorname{det}(P) P(\mathbf{t} \times \mathbf{n})=P \mathbf{b}$, so $\mathbf{B}=P \mathbf{b}$.
(To see the last fact, let $P$ be orthogonal and let $\mathbf{v}_{1}, \mathbf{v}_{2}, \mathbf{v}_{3}$ be the orthonormal columns of $P$. Then if $P$ is direct, we have $1=\operatorname{det}(P)=\left(\mathbf{v}_{1} \times \mathbf{v}_{2}\right) \cdot \mathbf{v}_{3}$, from the standard algorithm for computing the derivative by going down the third column. So in particular the unit vector $\mathbf{v}_{1} \times \mathbf{v}_{2}$ is $\mathbf{v}_{3}$, and the columns of $P$ form a righthanded orthonormal system. Then if $\mathbf{a}=\left(a_{1}, a_{2}, a_{3}\right)$ and $\mathbf{c}=\left(c_{1}, c_{2}, c_{3}\right)$, we can compute the cross-product $P \mathbf{a} \times P \mathbf{c}=\left(a_{1} \mathbf{v}_{1}+a_{2} \mathbf{v}_{2}+a_{3} \mathbf{v}_{3}\right) \times\left(b_{1} \mathbf{v}_{1}+b_{2} \mathbf{v}_{2}+b_{3} \mathbf{v}_{3}\right)=$ $\left(a_{1} b_{2}-b_{1} a_{2}\right) \mathbf{v}_{3}+\left(a_{3} b_{1}-b_{3} a_{1}\right) \mathbf{v}_{2}+\left(a_{2} b_{3}-b_{2} a_{3}\right) \mathbf{v}_{3}=P(\mathbf{a} \times \mathbf{b})$.
- (2.3.13) Recall that a curve is a generalized helix if it makes a fixed angle with some unit vector. Out curve is $\gamma(t)=\left(e^{\lambda t} \cos t, e^{\lambda t} \sin t, e^{\lambda t}\right)$, with tangent vector $\dot{\gamma}(t)=\left(\lambda e^{\lambda t} \cos t-e^{\lambda t} \sin t, \lambda e^{\lambda t} \sin t+e^{\lambda t} \cos t, \lambda e^{\lambda t}\right.$. Note that this vector has length $\sqrt{2 \lambda^{2}+1} e^{\lambda t}$. Now if $\mathbf{u}=(0,0,1)$, then $\dot{\gamma}(t) \cdot \mathbf{u}=\lambda e^{\lambda t}$, so if $\theta$ is the angle between $\dot{\gamma}(t)$ and $\mathbf{u}$, then $\cos \theta=\frac{\sqrt{2 \lambda^{2}+1}}{\lambda}$. We conclude $\theta$ is constant.

We also need to check the curvature of $\gamma$ is nonzero. But note that $\ddot{\gamma}(t)=\left(\left(\lambda^{2}-\right.\right.$ 1) $\left.e^{\lambda t} \cos t-2 \lambda e^{\lambda t} \sin t,\left(\lambda^{2}-1\right) e^{\lambda t} \sin t+2 \lambda e^{\lambda t} \cos t, \lambda^{2} e^{\lambda t}\right)$, so the first entry of $\ddot{\gamma} \times \dot{\gamma}$ is $-e^{\lambda t}\left(\lambda \sin t+3 \lambda^{2} \cos t\right)$ and the second is $\left(\lambda e^{\lambda t}\left(\lambda \cos t+3 \lambda^{2} \sin t\right)\right.$. These two terms cannot be zero simultaneously, so $\ddot{\gamma} \times \dot{\gamma}$ is not the zero vector for any $t$. Hence $\gamma$ has nonzero curvature anywhere.

